

## Propagation characteristic of ultrawide bandwidth Airy surface plasmons

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Grating





finite size grating — finite bandwidth

Grating



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Airy surface plasmon polaritons (Airy SPPs) propagating surface plasmons (SPPs) consisting of the properties of Airy beams.

Non diffracting beam remains intensity invariant during propagation.

Mathieu





**Basic characteristics** 

- Non diffracting beams contains infinite power or energy.
- Can be generated through conical superposition of plane waves



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M. Mazilu, D. J. Stevenson, F. Gunn-Moore, and K. Dholakia, *Laser Photonics Rev.* 4, 529 (2010). or J. Durnin, J. J. Miceli, and J. H. Eberly, PRL **58**, 1499 (1987). Check out to which these figures belong

Non diffracting beam remains intensity invariant during propagation.

Bessel

Mathieu



**Basic characteristics** 

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- Non diffracting beams contains infinite power or energy.
- Solution of 1D paraxial wave equation

M. Mazilu, D. J. Stevenson, F. Gunn-Moore, and K. Dholakia, Laser Photonics Rev. 4, 529 (2010).

G. A. Siviloglou and D. N. Christodoulides Opt. Lett. 32, 979 (2007).

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### Properties

- Airy beams can exist in a planar system, so particularly suitable for plasmonics.
- Accelerating
- Self bending
- Self healing

M. Mazilu, D. J. Stevenson, F. Gunn-Moore, and K. Dholakia, Laser Photonics Rev. 4, 529 (2010).

G. A. Siviloglou and D. N. Christodoulides Opt. Lett. 32, 979 (2007).

## Airy Plasmon

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A. Salandrino and D.N. Christodoulides, Opt.Lett. 35, 2082 (2010)

The magnetic field amplitude of the non-paraxial Airy plasmon in Fourier-space is given by

$$\tilde{A}_0(k_x; x_0, a) = \frac{1}{2\pi} \exp\left(-ak_x^2 x_0^2 + \frac{i}{3}[k_x^3 x_0^3 - 3a^2 k_x x_0 - ia^3]\right)$$
(1)

This yields the following magnetic field in real-space:

z [hm]

$$\mathbf{H}(x, y, z; x_0, a) = \hat{\mathbf{H}}(x, y; x_0, a) \begin{cases} \exp(-\alpha_I z) & \text{for } z > 0\\ \exp(+\alpha_{II} z) & \text{for } z < 0 \end{cases}$$
(2)

with  

$$\hat{\mathbf{H}}(x,y;x_{0},a) = \int_{-\infty}^{\infty} \tilde{A}_{0}(k_{x};x_{0},a) \left(\frac{q(k_{x})}{\beta}e_{x} - \frac{k_{x}}{\beta}e_{z}\right) \exp\left(ik_{x}x + iq(k_{x})z\right) dk_{x}$$
(3)
  
squared magnitude of normal E-field
$$\int_{-\infty}^{0} \frac{1}{20} \int_{0,1}^{0} \int_{0$$

## Simulation and experiment



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- Simulation: Lumerical FDTD solver
- Plane wave illumination (TFSF source)
- Angle of incidence: normal
- Spatial resolution 5 nm (uniform grid)

#### Advantage

- Computation region is divided into two distinct regions; total field and only scattered field
- suitable for non periodic structures in a multilayer system

- Experiment: Photoemission electron microscope (PEEM)
- Photoemitted yield are direct map of resultant field at the interface
- Photoemission is non-linear n-photon process
- Field of view is 80 µm (scale 660 pixels = 80µm)

electron counts





Results





Results





## Spatiotemporal evolution of Airy plasmon pulse

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- Source: 20fs Gaussian pulse
- Normal incidence
- Central wavelength 745nm,
- Bandwidth 40nm
- Pulse offset 50fs.
- temporal resolution 0.033fs

Time averaged Intensity of Airy plasmon pulse

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$$Intensity = \int \left[ |E_x(x, y, t)|^2 + |E_y(x, y, t)|^2 + |E_z(x, y, t)|^2 \right] dt$$



## Spatiotemporal evolution of Airy plasmon pulse

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- Source: 20fs Gaussian pulse
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temporal evolution of Airy bullet t=32.724fs 2.5 20 2 15 (und) z 1.5 1 5 0.5 0 0 -20 -15 -10 5 10 -5 0 x (µm)

Time averaged Intensity of Airy plasmon pulse

$$Intensity = \int \left[ |E_x(x, y, t)|^2 + |E_y(x, y, t)|^2 + |E_z(x, y, t)|^2 \right] dt$$



- Does not need diffraction and dispersion equalization
- possible under any dispersion conditions





# Thank You!!!

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